

Evidence for a stratified accretion disk wind in AGN

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Abstract. We present observational evidence supporting the presence of a stratified accretion disk wind in active galactic nuclei (AGN), based on multi-wavelength spectroscopic analysis of broad and narrow emission lines. The diversity in emission line profiles, ionization potentials, and kinematic signatures suggests a structured outflow emerging from the accretion disk, with different zones contributing to specific spectral features. High-ionization lines (e.g., CIV $\lambda 1549$) exhibit strong blueshifts and asymmetric profiles indicative of fast, inner winds, while low-ionization lines (e.g., H β , MgII $\lambda 2800$) show more symmetric profiles consistent with predominant emission from slower, denser regions farther out, although exhibiting systematic blueshifts in quasars radiating at high Eddington ratios. The intermediate ionization lines (e.g., AlIII $\lambda 1860$) present a situation that is intermediate in terms of shift amplitudes, although in several super-Eddington candidates radial outflow velocity may reach values comparable to the ones of the high ionization lines. These results are consistent with radiatively driven wind models featuring radial stratification. We made preliminary photoionization modeling assuming unabsorbed radiation emitted from the corona and the hotter disk regions emission or absorbed by a layer of gas. Our findings provide new constraints on the geometry and physical conditions of AGN winds, providing clear evidence in favor of stratified wind emission.

Key words: Supermassive black holes (1663) – Active galactic nuclei (16) – Quasars (1319) – Spectroscopy (1558) – Photoionization (2060)

1. Introduction

The concept of a quasar "main sequence" (MS) emerged from the landmark work of [Boroson & Green \(1992\)](#), who identified a strong anti-correlation between the relative strength of FeII $\lambda 4570$ emission and the full width at half

maximum (FWHM) of the broad $H\beta$ line. This relation, commonly expressed through the parameter $R_{\text{FeII}} = \text{FeII}\lambda4570/\text{H}\beta$, was first hinted at in earlier works (e.g., [Gaskell 1985](#) and has since been confirmed and extended in a large number of studies ([Sulentic et al., 2000a,b](#); [Shen & Ho, 2014](#); [Rakshit et al., 2020](#); [Wu & Shen, 2022](#)). The MS framework provides a powerful tool for organizing the remarkable diversity of type-1 active galactic nuclei (AGN) spectroscopic properties.

Within this scheme, type-1 AGN can be broadly separated into two main populations: Population A and Population B ([Sulentic et al., 2000a, 2002, 2011](#)). Population A sources, characterized by narrower $H\beta$ profiles (FWHM $< 4000 \text{ km s}^{-1}$), are generally associated with high accretion rates relative to the Eddington limit, while Population B sources (FWHM $> 4000 \text{ km s}^{-1}$) are typically lower Eddington ratio (L/L_{Edd} , where L_{Edd} is the Eddington luminosity) systems. A small subset, on the order of $\sim 10\%$ of the population, represents "extreme" Population A sources ($R_{\text{FeII}} > 1$), which are widely interpreted as candidates for super-Eddington accretion ([Wang et al., 2013](#); [Marziani & Sulentic, 2014](#); [Du et al., 2018](#); [Panda & Marziani, 2023](#); [Panda, 2024](#); [Marziani et al., 2025](#)).

The MS is thought to be primarily driven by the Eddington ratio L/L_{Edd} ([Marziani et al., 2001](#); [Boroson, 2002](#); [Shen & Ho, 2014](#); [Sun & Shen, 2015](#); [Panda et al., 2019](#)), with orientation effects also playing a significant role in shaping the observed diversity. Black hole mass effects, concomitant with viewing angle effects, become relevant when AGN samples cover a wide range in luminosity ([Marziani et al., 2018a](#); [Naddaf et al., 2025](#)). The organization of type-1 AGN properties across the MS has been extensively studied at multiple wavelengths, most prominently within the so-called 4D Eigenvector 1 (4DE1) parameter space introduced by [Sulentic et al. \(2000c\)](#). This framework combines optical, UV, and X-ray measures to capture the multidimensional diversity of quasars, while also providing a physical interpretation in terms of accretion physics and geometry (see [Fraix-Burnet et al. 2017](#) for a summary).

Large surveys, such as the Sloan Digital Sky Survey (SDSS, [York et al. 2000](#)), have provided further statistical foundation for these studies. For instance, [Zamfir et al. \(2010\)](#) analyzed ≈ 470 quasars at $z < 0.7$ with average bolometric luminosities $\log L \sim 45.5 \text{ erg/s}$, confirming the prevalence of the MS trends and strengthening the case for L/L_{Edd} as a primary driver. More recently, refinements in spectral analysis and improved databases have continued to reinforce this picture (e.g., [Shen & Ho 2014](#); [Wu & Shen 2022](#)).

In summary, the quasar main sequence organizes type-1 AGN along a continuum of properties primarily governed by Eddington ratio and orientation. At one end of the sequence lie the extreme Population A sources, with the largest FeII strengths and the highest inferred L/L_{Edd} values, plausibly representing systems accreting at or above the Eddington limit. These sources provide critical laboratories for understanding black hole growth under extreme conditions and the broader role of AGN feedback in galaxy evolution ([Marziani et al.,](#)

2025). Luminosity and black hole mass (M_{BH}) effects appear in sample covering large ranges in luminosity and mass. They can be reconduced to two effects: an increase in the amplitude of shifts with respect to the rest frame observed mainly in high-ionization lines (e.g., Marziani et al., 2016), and a prominent redward asymmetries associated with very massive black holes (Marziani et al., 2009; Marziani, 2023).

In this paper, we first present a brief summary of three recent works dealing with the observations of outflows as diagnosed from the shifts with respect to the rest frame, namely the prototypical $\text{CIV}\lambda 1549$ high-ionization line, the intermediate ionization emission $\text{AlIII}\lambda 1860$, and the low ionization line of $\text{MgII}\lambda 2800$ (Sections 2, 3). The three lines are all unresolved doublets associated with the resonance transition $^2P_{\frac{1}{2},\frac{3}{2}} \rightarrow ^2S_{\frac{1}{2}}$, with parent ionic species of different ionization potential, from ≈ 50 eV (CIV) to ≈ 25 eV (AlIII) and ≈ 15 eV (MgII). We afterwards attempt to explain the observational results by photoionization models focused on the sectors of the MS where the larger amplitude blueshifts are found, namely Population A and extreme Population A (Section 4).

2. $\text{CIV}\lambda 1549$ Shifts and outflows along the quasar main sequence

Systematic blueshifts of the high-ionization $\text{CIV}\lambda 1549$ emission line provide one of the clearest observational signatures of quasar outflows. Early evidence for such features in composite spectra of radio-quiet quasars was strengthened by HST/FOS studies (Marziani et al., 1996; Corbin & Boroson, 1996; Sulentic et al., 2007; Marziani et al., 2010; Leighly & Moore, 2004; Richards et al., 2011; Sulentic et al., 2007; Marziani et al., 2010; Sulentic et al., 2017). These analyses showed that CIV profiles can be decomposed into at least two components: a virialized, symmetric, and generally unshifted component (well represented by a Lorentzian in Population A sources), and a blue-shifted component that is naturally interpreted as emission from an outflowing wind. The latter is best modeled with a skewed Gaussian profile and becomes increasingly prominent at high R_{FeII} values, i.e., in extreme Population A quasars.

Population trends are striking at low z : large blueshifts of $v_r < -1000 \text{ km s}^{-1}$ are observed almost exclusively in Population A sources, particularly those with $\text{FWHM}(\text{H}\beta) < 4000 \text{ km s}^{-1}$ and strong FeII emission ($R_{\text{FeII}} > 1$; Zamfir et al. 2010; Marziani et al. 2010; Richards et al. 2011). In contrast, Population B quasars rarely show strong CIV blueshifts (if we exclude very high luminosity sources), reinforcing the link between high L/L_{Edd} and the presence of radiatively driven winds.

At higher luminosities, large blueshifts persist. Near-infrared spectroscopy of 52 Hamburg-ESO quasars with $1 \lesssim z \lesssim 3$ and $\log L > 47 \text{ erg/s}$ (HE main sequence, HEMS sample) confirmed that strong CIV blueshifts are associated with high Eddington ratios rather than simply high black hole masses or luminosities

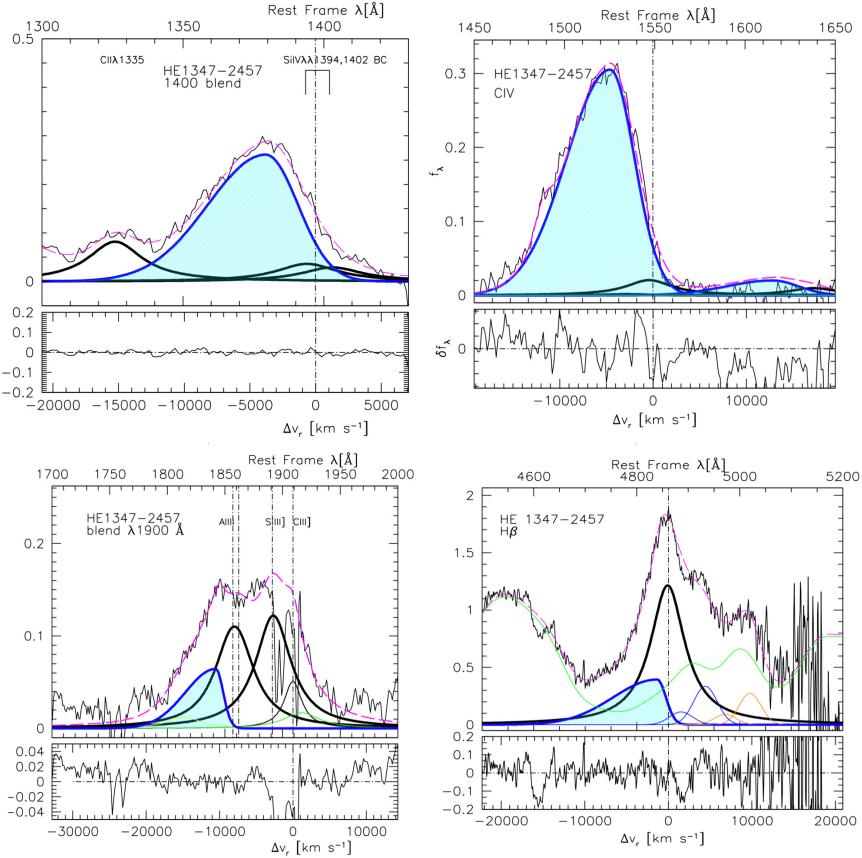


Figure 1. The continuum subtracted spectrum of the luminous Hamburg ESO Main Sequence (HEMS) quasar HE1347-2457, with an emphasis on the outflowing components of SiIVλ1397+OIV]λ1402, CIVλ1549, AlIIIλ1860, H β (cyan-shaded components).

(Marziani et al., 2018b; Richards et al., 2011; Giustini & Proga, 2019; Deconto-Machado et al., 2023, 2024). When viewed in the $L/L_{\text{Edd}}-L$ and $L/L_{\text{Edd}}-M_{\text{BH}}$ planes, blueshift amplitude is most clearly organized by accretion state, with Population A objects (above $L/L_{\text{Edd}} \approx 0.2$) being wind-dominated (Marziani et al., 2016).

There is, however, a trend between outflow velocity and luminosity (Marziani et al., 2016; Sulentic et al., 2017). The scaling of outflow velocity with luminosity further supports a radiative-driving mechanism. Weak but consistent correlations ($v \propto L^{0.25}$) are in line with predictions from radiation-driven disk-wind models (Murray et al., 1995; Laor & Brandt, 2002; Proga, 2007b,a). Pop. B at high luminosity also show large blueshifts, up to a few thousands km s $^{-1}$ (Sulentic

et al., 2017). There is evidence of Pop. B outflows at luminosity lower than $\log L \sim 47$ erg/s (e.g., Richards et al. 2011; Marziani et al. 2022a). Since Pop. B shifts at low luminosity remain of small amplitude and difficult to measure, the considerations presented in this paper are meant for Pop. A and extreme Pop. A only.

An example of powerful wind-dominated quasar of the HEMS survey is shown in Fig. 1: the profiles of high ionization lines such as CIV and SiIV $\lambda 1397 +$ OIV $\lambda 1402$ are dominated by blueshifted emission associated with the outflows. In this rather extreme cases, the outflows also significantly affect H β and the AlIII $\lambda 1860$ lines. While the bolometric luminosity sets the available radiative power, it is the Eddington ratio that regulates the efficiency of wind launching and explains the prevalence of large blueshifts in high- L/L_{Edd} sources.

Taken together, the restrictions of large CIV blueshifts to Population A and extreme Population A sources reinforces the interpretation of the quasar main sequence as primarily driven by L/L_{Edd} , with CIV outflows providing a direct tracer of radiative wind activity in the broad-line region.

3. Intermediate- and low-ionization lines: AlIII $\lambda 1860$ and MgII $\lambda 2800$

While the most dramatic blueshifts are observed in high-ionization lines such as CIV, intermediate- and low-ionization species also show systematic kinematic signatures that can provide complementary insights into quasar outflows.

3.1. AlIII $\lambda 1860$

The AlIII $\lambda 1860$ doublet, with intermediate ionization potential, generally shows only modest outflow signatures compared to CIV. Statistical analyses reveal that AlIII $\lambda 1860$ blueshifts are correlated with those of CIV but with a much shallower slope (≈ 0.1), indicating a less prominent wind contribution (Marziani et al., 2022b; Buendia-Rios et al., 2023). Quantitatively, the median centroid shift of AlIII $\lambda 1860$ follows the relation $c(\frac{1}{2})(\text{AlIII}) \approx (0.11 \pm 0.03) c(\frac{1}{2})(\text{CIV}) + (50 \pm 70) \text{ km s}^{-1}$ (Marziani et al., 2022b), consistent with a weaker kinematic response of the intermediate-ionization gas. However, in the most extreme Population A quasars ($R_{\text{FeII}} \gtrsim 1$), AlIII can display large blueshifts and asymmetric profiles comparable in strength to those of CIV, suggesting that under super-Eddington conditions the wind dominates across a broader ionization range.

3.2. MgII $\lambda 2800$

The MgII $\lambda 2800$ resonance doublet, one of the most widely used virial estimators of black hole mass (Marziani et al., 2013a), is typically far less affected by outflows (Trakhtenbrot & Netzer, 2012). Most quasars show symmetric and only weakly shifted MgII profiles. Nonetheless, careful spectral analysis has revealed

subtle displacements of the line core, along with FWHM increases of a few hundred km s^{-1} , in the highest L/L_{Edd} extreme Pop. A sources (Marziani et al., 2013b). These modest but systematic shifts indicate that even low-ionization gas can participate in winds under conditions of extreme accretion.

Following Popović et al. (2019), $\text{MgII}\lambda 2800$ shows Lorentzian MgII profiles and also signatures of outflow, but these are present only in part of the line and not in every source, as already noted by Marziani et al. (2013a). Popović et al. (2019) argue that MgII consists of two kinematic pieces": (i) a core that behaves similarly to $\text{H}\beta$ and appears virialized, and (ii) an additional "fountain-like" component, with motions roughly perpendicular to the disc, which produces the very broad wings and can reach shifts of a few thousand km s^{-1} . This second component may be associated with a photoionized "bowl" connecting the outer accretion disc and the inner torus (Goad et al., 2012), a configuration that is also consistent with a failed, radiatively accelerated, dusty wind (Czerny & Hrynowicz, 2011; Czerny et al., 2017; Naddaf et al., 2025). Since the fountain component does not correlate with the virial broadening estimators, it is interpreted as due to outflows/inflows rather than rotation.

This configuration can explain the Lorentzian wings of $\text{MgII}\lambda 2800$. It is, however, not intended to account for the net blueshift of the MgII line core reported by Marziani et al. (2013b) for the highest L/L_{Edd} sources. Those blueshifts are more naturally explained by a systematic line displacement, analogous to what is observed in $\text{CIV}\lambda 1549$, although with a smaller amplitude.

3.3. Ionization potential and the hierarchy of shifts

In extreme Pop. A quasars, $v(\text{AlIII}) \approx 0.3 v(\text{CIV})$ (and in some cases they are almost equal), while $v(\text{MgII}) \approx 0.1 v(\text{CIV})$ (Marziani et al., 2013b; Buendia-Rios et al., 2025). The comparative behavior of $\text{CIV}\lambda 1549$, $\text{AlIII}\lambda 1860$, and $\text{MgII}\lambda 2800$ therefore points to a simple trend: the higher the ionization potential of the line, the larger the blueshift. High-ionization lines such as CIV ($\text{IP} \approx 47.9 \text{ eV}$) show the largest shifts, often $> -1000 \text{ km s}^{-1}$ in extreme Pop. A sources. Intermediate-ionization AlIII ($\text{IP} \approx 28.4 \text{ eV}$) usually shows smaller shifts, but in the most extreme Pop. A quasars its outflow signature can approach that of CIV . In contrast, the low-ionization MgII doublet ($\text{IP} \approx 15 \text{ eV}$) shows only modest shifts, typically $\sim 10\%$ of the $\text{CIV}\lambda 1549$ amplitude (Marziani et al., 2013b, 2022b; Buendia-Rios et al., 2023, 2025). This ionization-stratified sequence matches the expectations of disk–wind models, in which the highest-ionization gas traces the fastest parts of the flow, while lower-ionization lines form deeper in the broad-line region, where motions are more nearly virial.

This stratification follows naturally if the high-ionization gas occupies the more accelerated layers of a radiatively driven disk wind close to the continuum source, whereas lower-ionization species arise farther out in the BLR (Murray et al., 1995; Proga, 2007b). In this picture, extreme Population A quasars lie in the regime where radiation pressure not only controls the dynamics of the

highest-ionization gas but also has a measurable impact on intermediate- and even low-ionization lines.

4. Spectral energy distributions of high accretors and photoionization simulations

The interpretation of outflow signatures in high-accretion quasars requires a realistic description of their spectral energy distributions (SEDs). Recent works have established that quasars with high Eddington ratios ($L/L_{\text{Edd}} \gtrsim 1$) and extreme FeII emission ($R_{\text{FeII}} \gtrsim 1$) exhibit consistent and distinctive SEDs (Jin et al., 2012; Marziani & Sulentic, 2014; Ferland et al., 2020; Panda & Marziani, 2023; Garnica et al., 2025). We utilized the median SED derived from a sample of ≈ 150 low-redshift Pop. xA quasars (Garnica et al., 2025) as well as the high L/L_{Edd} SED from Jin et al. (2012), and we assume that the line emitting gas is exposed either to the unobscured continuum or continuum absorbed by a hot, dense layer of gas at $r \approx 200r_g$ (for $\approx 10^8 M_{\odot}$; $\log N_{\text{H}} = 10^{23} \text{ cm}^{-2}$; $\log n_{\text{H}} = 10^{11} \text{ cm}^{-3}$). A lower black hole mass $10^7 M_{\odot}$ appropriate for NLSy1s would require a different, “harder” SED.

4.1. Radiatively driven outflows

Photoionization simulations highlight the conditions under which radiatively driven winds can be launched. Outflows are possible at relatively low column densities ($N_{\text{c}} \approx 10^{21} \text{ cm}^{-2}$), where radiative acceleration exceeds gravity. The requirement that the force multiplier $M > 2$ translates into the condition $M(L/L_{\text{Edd}}) > 1$ (see also Ferland et al., 2009; Netzer & Marziani, 2010). At higher column densities ($N_{\text{c}} \approx 10^{24} \text{ cm}^{-2}$; Ferland & Persson 1989; Panda et al. 2020; Panda 2021), radiation fails to overcome gravity, and the gas remains bound. These results confirm that outflow efficiency is intimately tied to the accretion state and gas structure in the inner broad-line region.

4.2. Predicted emission luminosities and line ratios

The simulations further predict the radial stratification of outflow emission. CIV $\lambda 1549$ emission is favored at smaller radii, AlIII $\lambda 1860$ at intermediate radii, and MgII at larger distances, reflecting both ionization potential and gas density conditions. The luminosities ($L(\text{CIV})$, $L(\text{AlIII})$, $L(\text{MgII})$) define a narrow “corridor” of optimal ionization parameter where strong emission can be sustained (Fig. 2). The continuum luminosity at $\lambda = 3000 \text{ \AA}$ has been assumed to be $\lambda L_{\lambda}(3000 \text{ \AA}) \approx 4.2 \cdot 10^{44} \text{ erg/s}$. The hatched parts identify areas of the parameter plane where radiation forces are insufficient to drive an outflow; dark blue parts are regions where line emission is exceedingly low. The case shown here refers to a partially-absorbed high L/L_{Edd} Jin et al. (2012) continuum, with the absorber located at $r \approx 200r_g$. Due to the relative proximity of the

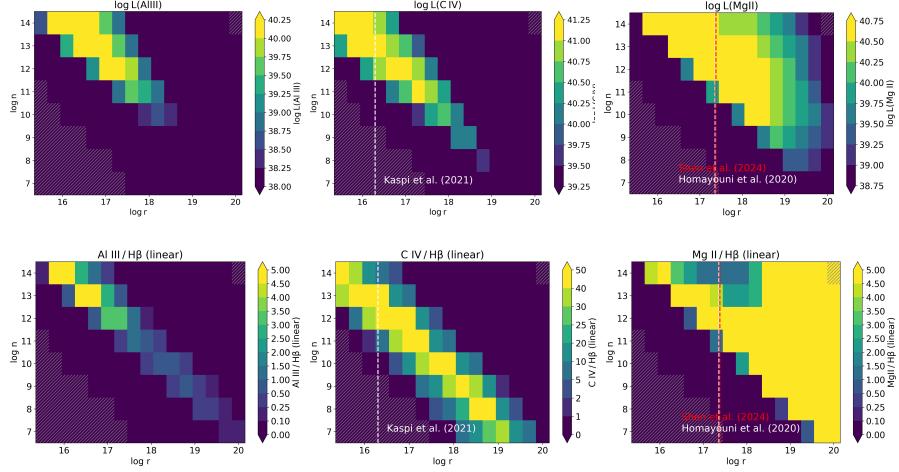


Figure 2. Behavior of line luminosity (top row) and intensity ratios with $H\beta$ (bottom), for $CIV\lambda 1549$, $AlIII\lambda 1860$, and $MgII\lambda 2800$ as a function of hydrogen density and radius, for a 10^8 solar masses black hole. The white vertical lines identify the radii measured according to the scaling laws of [Kaspi et al. \(2021\)](#) for $CIV\lambda 1549$ and [Homayouni et al. \(2020\)](#); [Shen et al. \(2024\)](#) for $MgII\lambda 2800$.

absorber relative to the central continuum source, the absorbing gas remains hot and producing significant absorption only in the range between 1 and 10 keV. A fully unabsorbed [Jin et al. \(2012\)](#) continuum would give rise to similar trends.

The corridor peak, if considered at fixed density, is displaced toward larger radii for $AlIII$, and even more so for $MgII$ with respect to CIV (top row of Fig. 2), as schematically emphasized in Fig. 3. Outside the corridor for CIV and $AlIII$ emission, for radii smaller than the ones of the corridor, over-ionization suppresses line production, a behavior consistent with the observed weakness of $CIV\lambda 1549$ in extreme Pop. A quasars ([Kaspi et al., 2021](#)). For larger radii, under-ionization depresses the emission of both $CIV\lambda 1549$ and $AlIII\lambda 1860$. The behavior of $MgII\lambda 2800$ is different: strong emission and large $MgII\lambda 2800/H\beta$ radii are possible over a wide range of radii that are however larger than the ones where $AlIII\lambda 1860$ and $CIV\lambda 1549$ are maximized.

Predicted line ratios also provide critical tests (bottom panels of Fig. 2). For instance, CLOUDY simulations reproduce observed $CIV\lambda 1549/H\beta$ ratios $> 10-20$

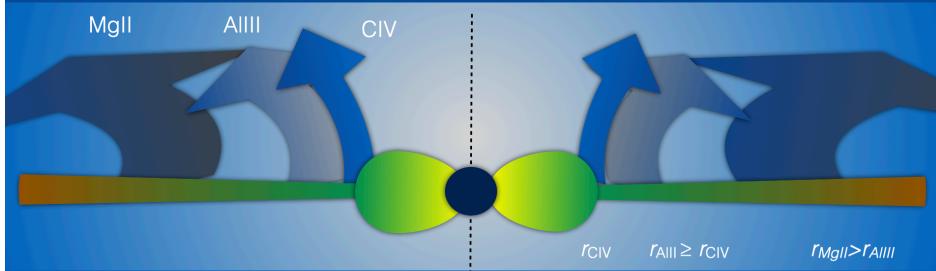


Figure 3. Sketch illustrating the differences in launching radii for the radiatively driven-wind emitting $\text{CIV}\lambda 1549$, $\text{AlIII}\lambda 1860$, $\text{MgII}\lambda 2800$. The accretion disk around the central black hole is assumed to have an inner, puffed-up and optically thick region sustained by radiation pressure, as well as an outer thin, optically thick region. The various elements are not drawn to scale.

only under specific combinations of density and ionization, while $\text{MgII}\lambda 2800/\text{H}\beta$ ratios match those measured in large quasar samples (Homayouni et al., 2022; Prince et al., 2023). The photoionization model therefore not only explains the ionization-dependent hierarchy of blueshifts ($\text{CIV}\lambda 1549 > \text{AlIII}\lambda 1860 > \text{MgII}\lambda 2800$) but also constrains the physical conditions under which these outflows form.

5. Conclusions

Our analysis of emission-line shifts along the quasar main sequence reinforces the view that radiatively driven outflows are a ubiquitous feature of type-1 AGN. Outflow signatures are present even in single-epoch spectra, and they become dominant in high-ionization lines at high Eddington ratios and, more generally, at the highest luminosities ($L \gtrsim 10^{47} \text{ erg s}^{-1}$).

In terms of accretion mode, Pop. A AGN satisfying the criterion $L/L_{\text{Edd}} > 0.1 - 0.2$ could be defined as black hole with an inner optically thick, geometrically thick region (c.f., Giustini & Proga, 2019). Among Population A sources, the amplitude of line blueshifts decreases systematically with ionization potential, from $\text{CIV}\lambda 1549$ (C^{3+}) to $\text{AlIII}\lambda 1860$ and down to $\text{MgII}\lambda 2800$. While $\text{CIV}\lambda 1549$ blueshifts can exceed several thousand km s^{-1} , $\text{MgII}\lambda 2800$ shows only subtle but measurable shifts, detectable mainly in extreme Pop. A quasars.

These trends are consistent with an ionization-stratified broad-line region, where higher ionization lines originate closer to the black hole in regions more directly exposed to radiation pressure (Figure 3).

Photoionization simulations with realistic high-accretor SEDs demonstrate that outflows can be launched over a wide range of radii if the column density remains low ($N_c \approx 10^{21} \text{ cm}^{-2}$), whereas dense gas ($N_c \approx 10^{24} \text{ cm}^{-2}$) remains gravitationally bound. These models reproduce both the hierarchy of shift amplitudes ($\text{CIV} > \text{AlIII} > \text{MgII}$) and the reverberation-mapped radii of $\text{CIV}\lambda 1549$ and $\text{MgII}\lambda 2800$ (Homayouni et al., 2020; Khadka et al., 2021; Cao et al., 2022; Shen et al., 2024), providing a physically consistent picture that links SED shape, BLR stratification, and wind dynamics.

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